

# Fekete-Szegő Functional and Second Hankel Determinant for a Subclass of Te-Univalent Functions Connected to Horadam Polynomial.

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**Abstract:** This paper presents a subclass, referred to as  $\mathfrak{X}_\sigma^{\alpha, \beta}[\mu_1; \nu_1, \Phi(x)]$  of Te-univalent function that employs Horadam polynomial. The research explores this subclass and establishes preliminary coefficient bounds for  $|b_2|, |b_3|, |b_4|$ , Fekete-Szegő inequality and second hankel determinant for this category. Furthermore several corollaries are included to clarify the significance of the results.

**KeyWords:** Analytic functions, Te-Univalent functions, Coefficient bounds, Horadam Polynomial and Fekete-Szegő Inequality.

## 1 Introduction

Let  $\psi(\mathfrak{z})$  be normalized analytic function expressed as follows

$$\psi(\mathfrak{z}) = \mathfrak{z} + b_2\mathfrak{z}^2 + b_3\mathfrak{z}^3 + \dots \quad (1.1)$$

where  $\mathfrak{z} \in \mathfrak{D}$  and  $\mathfrak{D} = \{\mathfrak{z} : \mathfrak{z} \in \mathbb{C} \text{ and } |\mathfrak{z}| < 1\}$ .

We denote the class of all such functions as  $\mathcal{A}$  and we focus on the subclass  $\mathcal{S}$  of  $\mathcal{A}$ , which is characterized by  $\mathcal{S} = \{\psi \in \mathcal{A} : \psi \text{ is univalent in } \mathfrak{D}\}$ .

It is well-known result that any function  $\psi$  that belongs to  $\mathcal{S}$  possesses an inverse function  $\psi^{-1}$  which can be represented in the following manner:

$$\psi^{-1}(w) = w + \sum_{n=2}^{\infty} d_n w^n. \quad (1.2)$$

such that  $\psi^{-1}(\psi(\mathfrak{z})) = \mathfrak{z}$  for  $\mathfrak{z} \in \mathfrak{D}$  and  $\psi(\psi^{-1}(w)) = w$  for  $w \in \mathfrak{D}$ , with  $|w| < \rho_0(\psi)$  and  $\rho_0(\psi) \geq \frac{1}{4}$ . Furthermore, the inverse function  $\psi^{-1}$ , as given in (1.2) can be rewritten as

$$\mathfrak{F}(w) = \psi^{-1}(w) = w - b_2 w^2 + (2b_2^2 - b_3)w^3 - (5b_2^3 - 5b_2 b_3 + b_4)w^4 + \dots \quad (w \in \mathfrak{D}) \quad (1.3)$$

The bi-univalent functions introduced by Lewin [13], which are defined as analytic function  $\psi$  within the unit disc  $\mathfrak{D}$ , where both  $\psi$  and its inverse  $\psi^{-1}$  are univalent in  $\mathfrak{D}$ . This category of functions is symbolized by  $\sigma$ . Examples of functions that are classified in  $\sigma$  include

$\psi_1(\mathfrak{z}) = \frac{\mathfrak{z}}{1-\mathfrak{z}}$ ,  $\psi_2(\mathfrak{z}) = -\log(1-\mathfrak{z})$ ,  $\psi_3(\mathfrak{z}) = \frac{1}{2} \log\left(\frac{1+\mathfrak{z}}{1-\mathfrak{z}}\right)$ . However the function  $\frac{\mathfrak{z}}{(1-\mathfrak{z})^2}$  is included in the class  $\mathcal{S}$  but not belong to  $\sigma$ .

For integers  $n \geq 1$  and  $q \geq 1$ , the  $q^{\text{th}}$  Hankel determinant given by

$$H_q(n) = \begin{vmatrix} b_n & b_{n+1} & \dots & b_{n+q-1} \\ b_{n+1} & b_{n+2} & \dots & b_{n+q} \\ \vdots & \vdots & \ddots & \vdots \\ b_{n+q-1} & b_{n+q} & \dots & b_{n+2q-1} \end{vmatrix} \quad (b_1 = 1)$$

The properties of Hankel determinant can be studied in [18].

Let the operator  $\mathfrak{I}: \mathcal{A} \rightarrow \mathcal{A}$  defined by the equation

$$\mathfrak{I}(\psi(z)) = z + \sum_{n=2}^{\infty} t_n b_n z^n \quad (1.4)$$

where  $z \in \mathcal{D}$  and  $t_n \in \mathbb{C}$ .

The concept of Te-univalence associated with an operator was established by Abd-Eltawab[1] which is a extension and generalization of the idea of bi-univalence of the function class defined in (1.4) on  $\mathcal{D}$ . Let  $\mathfrak{I}_\sigma$  be the class of all functions given by (1.4) that are univalent in  $\mathcal{D}$ . Every  $\mathfrak{I}\psi \in \mathfrak{I}$  has an inverse defined as  $(\mathfrak{I}\psi)^{-1}$ .

$$(\mathfrak{I}\psi)^{-1}((\mathfrak{I}\psi)(z)) = z$$

and

$$\mathfrak{I}((\mathfrak{I}\psi)^{-1}(w)) = w \quad \left( |w| < \rho_0(\mathfrak{I}\psi); \rho_0(\mathfrak{I}\psi) \geq \frac{1}{4} \right)$$

where

$$(\mathfrak{I}\psi)^{-1}(w) = w - t_2 b_2 w^2 + (2t_2^2 b_2^2 - t_3 b_3) w^3 - (5t_2^3 b_2^3 - 5t_2 t_3 b_2 b_3 + t_4 b_4) w^4 + \dots$$

If both  $\mathfrak{I}\psi$  and  $(\mathfrak{I}\psi)^{-1}$  are univalent in  $\mathcal{D}$ , the function  $\psi$  given by (1.1) is Te-univalent in  $\mathcal{D}$  associated with  $\mathfrak{I}$ . Let  $\mathfrak{I}_\sigma$  be the class of all function given by (1.1) that are Te-univalent in  $\mathcal{D}$  and associated with  $\mathfrak{I}$

Note that for  $\mathfrak{I}\psi = \psi$  we have  $\mathfrak{I}_\sigma = \sigma$  and if  $t_n \neq 1$  for some  $n$ , then

$$\mathfrak{I}\psi(\mathfrak{I}\mathfrak{I}(w)) = w + 2(t_3 - t_2^2) b_2^2 w^3 + \dots \neq w$$

For function  $\psi$  given by (1.1) and  $\chi$  is given by

$$\chi(z) = z + \sum_{n=2}^{\infty} c_n z^n.$$

The Hadamard product of  $\psi$  and  $\chi$  is defined by

$$(\psi * \chi)z = z + \sum_{n=2}^{\infty} b_n c_n z^n = (\chi * \psi)z.$$

For complex parameters  $\mu_1, \mu_2, \mu_3, \dots, \mu_q$  and  $\nu_1, \nu_2, \nu_3, \dots, \nu_\zeta$  ( $\nu_j \neq 0, -1, -2, \dots, j = 1, 2, \dots, \zeta$ ) the generalized hypergeometric function  ${}_q\mathfrak{I}_\zeta$  is defined as follows

$${}_q\mathfrak{I}_\zeta(\mu_1, \mu_2, \mu_3, \dots, \mu_q; \nu_1, \nu_2, \nu_3, \dots, \nu_\zeta; z) = \sum_{n=0}^{\infty} \frac{(\mu_1)_n \dots (\mu_q)_n z^n}{(\nu_1)_n \dots (\nu_\zeta)_n n!}, \quad (z \in \mathcal{D})$$

where  $q, \zeta \in N_0 := N \cup \{0\}$  with  $\zeta + 1 \geq q$  and  $(\gamma)_n$  is the shift factorial(or Pochhammer symbol) defined in terms of the gamma function  $\Gamma$ , by

$$(\gamma)_n = \frac{\Gamma(\gamma+n)}{\Gamma(\gamma)} = \begin{cases} 1, & \text{if } n = 0 \\ \gamma(\gamma + 1) \dots (\gamma + n - 1), & \text{for } n \in N \end{cases}$$

corresponding a function

$$\mathfrak{I}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta; z) = z {}_q\mathfrak{I}_\zeta(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta; z) \quad (z \in \mathcal{D}) \quad (1.5)$$

Dziok and Srivastava [6] investigated that a linear operator is defined as

$$\mathfrak{G}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta): \mathcal{A} \rightarrow \mathcal{A}$$

defined as the following: Hadamard Product

$$\mathfrak{G}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta)\psi(z) = \mathfrak{I}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta; z) * \psi(z)$$

where  $\mathfrak{I}$  is defined as (1.5), and  $q, \zeta \in N_0$  with  $\zeta + 1 \geq q$ .

If  $\psi \in \mathcal{A}$  is given by (1.1) then

$$\mathfrak{G}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta)\psi(z) = z + \sum_{n=2}^{\infty} \Gamma_n[\mu_1; \nu_1] b_n z^n \quad (z \in \mathcal{D}) \quad (1.6)$$

where

$$\Gamma_n[\mu_1; \nu_1] = \frac{(\mu_1)_{n-1} \dots (\mu_q)_{n-1}}{(\nu_1)_{n-1} \dots (\nu_\zeta)_{n-1}} \frac{1}{(n-1)!} \quad (1.7)$$

To simplify the notation, we write:

$$\mathfrak{G}(\mu_1, \mu_2, \dots, \mu_q; \nu_1, \nu_2, \dots, \nu_\zeta)\psi = \mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi$$

Let  $\mathfrak{T}_\zeta^{q,\zeta}$  represents the class of functions given by (1.6) that are univalent in  $\mathfrak{D}$ . Every function  $\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi \in \mathfrak{T}_\zeta^{q,\zeta}[\mu_1; \nu_1]$  has an inverse,  $g = (\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi)^{-1}$  defined as [14]

$$g(\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi(\mathfrak{z})) = \mathfrak{z} \quad (\mathfrak{z} \in \mathfrak{D}),$$

and

$$\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi(g(w)) = w, \quad (|w| < \rho_0(\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi) \text{ and } \rho_0(\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi) \geq \frac{1}{4})$$

where

$$\begin{aligned} g(w) &= (\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi)^{-1}(w) \\ &= w - \Gamma_2[\mu_1; \nu_1]b_2w^2 + [2(\Gamma_2[\mu_1; \nu_1])^2b_2^2 - \Gamma_3[\mu_1; \nu_1]b_3]w^3 \\ &\quad - [5(\Gamma_2[\mu_1; \nu_1])^3b_2^3 - 5\Gamma_2[\mu_1; \nu_1]\Gamma_3[\mu_1; \nu_1]b_2b_3 + \Gamma_4[\mu_1; \nu_1]b_4]w^4 + \dots \end{aligned} \quad (1.8)$$

and  $\Gamma_n[\mu_1; \nu_1]$  is given by (1.7). A function  $\psi$  given by (1.1) is said to be Te-univalent [1] in  $\mathfrak{D}$  associated with operator  $\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]$ . If both  $\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi$  and  $(\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi)^{-1}$  are univalent in  $\mathfrak{D}$ . Let  $\mathfrak{T}_\sigma^{q,\zeta}[\mu_1; \nu_1]$  be the class of all functions given by (1.1) that are Te-univalent in  $\mathfrak{D}$  associated with  $\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]$ .

Note that for  $q = 2, \zeta = 1, \mu_1 = \nu_1 = c$  and  $\mu_2 = 1$  we have  $\mathfrak{T}_\sigma^{q,\zeta}[c, 1; c] = \sigma$  and if  $\Gamma_n[\mu_1; \nu_1] \neq 1$  for some  $n$ , we have

$$\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\psi(\mathfrak{G}_{q,\zeta}[\mu_1; \nu_1]\mathfrak{F}(w)) = w + 2[\Gamma_3[\mu_1; \nu_1] - (\Gamma_2[\mu_1; \nu_1])^2]b_2^2w^3 + \dots \neq w.$$

where  $\mathfrak{F}$  is given by (1.3).

In 1985, Horadam and Mohan [9] defined the Horadam polynomials  $\varphi_n(x) = \varphi_n(\gamma, \delta; l, m)$  by the following recurrence relation:

$$\varphi_n(x) = l x \varphi_{n-1}(x) + m \varphi_{n-2}(x), \text{ for } n \geq 3. \quad (1.9)$$

with initial values

$$\varphi_1(x) = \gamma, \varphi_2(x) = \delta x \text{ and } \varphi_3(x) = l\delta x^2 + m\gamma. \quad (1.10)$$

Further, the generating function of Horadam polynomial is

$$\Phi(x) = \sum_{n=1}^{\infty} \varphi_n(x) \mathfrak{z}^{n-1} = \frac{\gamma + (\delta - \gamma l)x \mathfrak{z}}{1 - l x \mathfrak{z} - m \mathfrak{z}^2}.$$

In this paper, the argument of  $x \in \mathbb{R}$  is independent of the argument  $\mathfrak{z} \in \mathfrak{D}$ ; that is  $x \neq \mathcal{R}(\mathfrak{z})$ . By giving values for  $\gamma, \delta, l$  and  $m$  the Horadam polynomials leads to several known polynomials as follows:

- If  $\gamma = \delta = l = m = 1$ , we get Fibonacci polynomials  $F_n(x)$
- If  $\gamma = 2$  and  $\delta = l = m = 1$ , we get Lucas polynomials  $L_n(x)$ .
- If  $\gamma = m = 1$  and  $\delta = l = 2$ , we get Pell polynomials  $P_n(x)$ .
- If  $\gamma = \delta = l = 2$  and  $m = 1$ , we get Pell-Lucas polynomials  $Q_n(x)$ .
- If  $\gamma = \delta = -m = 1$  and  $l = 2$ , we get the Chebyshev Polynomials  $T_n(x)$  of the first kind.
- If  $\gamma = -m = 1$  and  $\delta = l = 2$  we get the Chebyshev Polynomials  $U_n(x)$  of the second kind.

For more information see [2, 3, 4, 5, 9, 10, 12, 15, 16, 17, 19].

## 2 Basic Preliminaries and Lemmas

**Definition:** A function  $\psi \in \sigma$  is said to be in the class  $\mathfrak{F}_{\sigma}^{q,s}[\mu_1; \nu_1, \Phi(x)]$  if it satisfies :

$$(1 - \rho) \frac{\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi(\mathfrak{z})}{\mathfrak{z}} + \rho (\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi(\mathfrak{z}))' < \Phi(x, \mathfrak{z}) + 1 - \gamma \quad (\mathfrak{z} \in \mathfrak{D}) \quad (2.1)$$

and

$$(1 - \rho) \frac{g(w)}{w} + \rho g'(w) < \Phi(x, w) + 1 - \gamma \quad (w \in \mathfrak{D}) \quad (2.2)$$

where  $g(w) = (\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi)^{-1}(w)$  ,  $<$  means subordination and  $0 \leq \rho \leq 1$ .

**Lemma 2.1** [11] Let  $k, j \in \mathbb{R}$  and  $x, y \in \mathbb{C}$ . If  $|x| < r$  and  $|y| < r$ ,

$$|(k + j)x + (k - j)y| \leq \begin{cases} 2|k|r, & \text{if } |k| \geq j \\ 2|j|r, & \text{if } |k| \leq j \end{cases}$$

**Lemma 2.2** [7] Suppose that  $\mathcal{P}$  is the set of all analytic functions of the form

$$p(\mathfrak{z}) = 1 + \sum_{n=1}^{\infty} \alpha_n \mathfrak{z}^n$$

satisfying  $R(p(\mathfrak{z})) > 0, \mathfrak{z} \in \mathbb{C}$  and  $p(0) = 1$ . Then  $|\alpha_n| \leq 2 \quad n = 1, 2, 3, \dots$  For any value of  $n = 1, 2, 3, \dots$ , this inequality is sharp. For example the function  $\alpha(\mathfrak{z}) = \frac{1+\mathfrak{z}}{1-\mathfrak{z}}$  is equal for all  $n$ .

**Lemma 2.3** [8] Suppose that  $\mathcal{P}$  is the set of all analytic functions of the form  $p(\mathfrak{z}) = 1 + \sum_{n=1}^{\infty} \alpha_n \mathfrak{z}^n$  satisfying  $R(p(\mathfrak{z})) > 0, \mathfrak{z} \in \mathbb{C}$  and  $p(0) = 1$ . Then

$$2\alpha_2 = \alpha_1^2 + (4 - \alpha_1^2)y$$

$$4\alpha_3 = \alpha_1^3 + 2(4 - \alpha_1^2)\alpha_1 y - (4 - \alpha_1^2)\alpha_1 y^2 + 2(4 - \alpha_1^2)(1 - |y|^2)\mathfrak{z}$$

for some  $y, \mathfrak{z}$  with  $|y| \leq 1, |\mathfrak{z}| \leq 1$ .

## 3 Coefficient Bounds

**Theorem 3.1** If  $\psi$  as defined in (1.1) is in  $\mathfrak{F}_{\sigma}^{q,s}[\mu_1; \nu_1, \Phi(x)]$ . Then

$$|b_2| \leq \frac{|\delta x| \sqrt{|\delta x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|[(2\rho+1)(\delta x+1)+\rho^2]\delta x - [l\delta x^2+m\gamma](\rho+1)^2|}} \quad (3.1)$$

$$|b_3| \leq \frac{|\delta x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{\delta x}{(\rho+1)^2} \right|. \quad (3.2)$$

and

$$|b_4| \leq \frac{|l^2 \delta x^3 + l m x \gamma + m x \delta|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{5x^2 \delta^2}{2|(\rho+1)(2\rho+1)|\Gamma_4[\mu_1; \nu_1]} \quad (3.3)$$

*Proof.* Let  $\psi \in \mathfrak{F}_{\sigma}^{q,s}[\mu_1; \nu_1, \Phi(x)]$ . Then

$$(1 - \rho) \frac{\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi(\mathfrak{z})}{\mathfrak{z}} + \rho (\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi(\mathfrak{z}))' = \Phi(x, c(\mathfrak{z})) + 1 - \gamma \quad (\mathfrak{z} \in \mathfrak{D}) \quad (3.4)$$

and

$$(1 - \rho) \frac{g(w)}{w} + \rho g'(w) = \Phi(x, d(w)) + 1 - \gamma \quad (w \in \mathfrak{D}) \quad (3.5)$$

where  $g(w) = (\mathfrak{G}_{q,s}[\mu_1; \nu_1] \psi)^{-1}(w)$  is given by (1.8).

Where  $\alpha, \beta \in \mathcal{P}$  as follows:

$$\alpha(\mathfrak{z}) = \frac{1+c(\mathfrak{z})}{1-c(\mathfrak{z})} = 1 + \alpha_1\mathfrak{z} + \alpha_2\mathfrak{z}^2 + \alpha_3\mathfrak{z}^3 + \dots$$

$$\Rightarrow c(\mathfrak{z}) = \frac{\alpha(\mathfrak{z})-1}{\alpha(\mathfrak{z})+1} \quad (\mathfrak{z} \in \mathfrak{D}) \quad (3.6)$$

and

$$\beta(w) = \frac{1+d(w)}{1-d(w)} = 1 + \beta_1w + \beta_2w^2 + \beta_3w^3 + \dots$$

$$\Rightarrow d(w) = \frac{\beta(w)-1}{\beta(w)+1} \quad (w \in \mathfrak{D}) \quad (3.7)$$

From (3.6)and (3.7)

$$c(\mathfrak{z}) = \frac{\alpha_1}{2}\mathfrak{z} + \left(\frac{\alpha_2}{2} - \frac{\alpha_1^2}{4}\right)\mathfrak{z}^2 + \left(\frac{\alpha_3}{2} - \frac{\alpha_1\alpha_2}{2} + \frac{\alpha_1^3}{8}\right)\mathfrak{z}^3 + \dots \quad (3.8)$$

and

$$d(w) = \frac{\beta_1}{2}w + \left(\frac{\beta_2}{2} - \frac{\beta_1^2}{4}\right)w^2 + \left(\frac{\beta_3}{2} - \frac{\beta_1\beta_2}{2} + \frac{\beta_1^3}{8}\right)w^3 + \dots \quad (3.9)$$

From (3.8) and (3.9)

$$\Phi(x, c(\mathfrak{z})) + 1 - \gamma = 1 + \frac{\varphi_2(x)}{2}\alpha_1\mathfrak{z} + \left[\frac{\varphi_2(x)}{2}\left(\alpha_2 - \frac{\alpha_1^2}{2}\right) + \frac{\varphi_3(x)}{4}\alpha_1^2\right]\mathfrak{z}^2 + \dots \quad (3.10)$$

and

$$\Phi(x, d(w)) + 1 - \gamma = 1 + \frac{\varphi_2(x)}{2}\beta_1w + \left[\frac{\varphi_2(x)}{2}\left(\beta_2 - \frac{\beta_1^2}{2}\right) + \frac{\varphi_3(x)}{4}\beta_1^2\right]w^2 + \dots \quad (3.11)$$

Also

$$(1 - \rho) \frac{\mathfrak{G}_{q,\varsigma}[\mu_1; \nu_1]\psi(\mathfrak{z})}{\mathfrak{z}} + \rho(\mathfrak{G}_{q,\varsigma}[\mu_1; \nu_1]\psi(\mathfrak{z}))' = 1 + (\rho + 1)\Gamma_2[\mu_1; \nu_1]b_2\mathfrak{z}$$

$$+ (2\rho + 1)\Gamma_3[\mu_1; \nu_1]b_3\mathfrak{z}^2 + (3\rho + 1)\Gamma_4[\mu_1; \nu_1]b_4\mathfrak{z}^3 + \dots \quad (3.12)$$

and

$$(1 - \rho) \frac{g(w)}{w} + \rho g'(w) = 1 - (\rho + 1)\Gamma_2[\mu_1; \nu_1]b_2w + (2\rho + 1)(2(\Gamma_2[\mu_1; \nu_1])^2b_2^2 - \Gamma_3[\mu_1; \nu_1]b_3)w^2$$

$$- (3\rho + 1)(5(\Gamma_2[\mu_1; \nu_1])^3b_2^3 - 5\Gamma_2[\mu_1; \nu_1]\Gamma_3[\mu_1; \nu_1]b_2b_3 + \Gamma_4[\mu_1; \nu_1]b_4)w^3 + \dots \quad (3.13)$$

From (3.10), (3.11), (3.12) and (3.13)

$$(\rho + 1)\Gamma_2[\mu_1; \nu_1]b_2 = \frac{\varphi_2(x)}{2}\alpha_1. \quad (3.14)$$

$$(2\rho + 1)\Gamma_3[\mu_1; \nu_1]b_3 = \frac{\varphi_2(x)}{2}\left(\alpha_2 - \frac{\alpha_1^2}{2}\right) + \frac{\varphi_3(x)}{4}\alpha_1^2. \quad (3.15)$$

$$(3\rho + 1)\Gamma_4[\mu_1; \nu_1]b_4 = \frac{\varphi_2(x)}{2}\left(\alpha_3 - \alpha_1\alpha_2 + \frac{\alpha_1^3}{4}\right) + \frac{\varphi_3(x)}{2}\alpha_1\left(\alpha_2 - \frac{\alpha_1^2}{2}\right) + \frac{\varphi_4(x)}{8}\alpha_1^3. \quad (3.16)$$

and

$$-(\rho + 1)\Gamma_2[\mu_1; \nu_1]b_2 = \frac{\varphi_2(x)}{2}\beta_1. \quad (3.17)$$

$$(2\rho + 1)(2(\Gamma_2[\mu_1; \nu_1])^2b_2^2 - \Gamma_3[\mu_1; \nu_1]b_3) = \frac{\varphi_2(x)}{2}\left(\beta_2 - \frac{\beta_1^2}{2}\right) + \frac{\varphi_3(x)}{4}\beta_1^2. \quad (3.18)$$

$$-(3\rho + 1)(5(\Gamma_2[\mu_1; \nu_1])^3b_2^3 - 5\Gamma_2[\mu_1; \nu_1]\Gamma_3[\mu_1; \nu_1]b_2b_3 + \Gamma_4[\mu_1; \nu_1]b_4) =$$

$$\frac{\varphi_2(x)}{2}\left(\beta_3 - \beta_1\beta_2 + \frac{\beta_1^3}{4}\right) + \frac{\varphi_3(x)}{2}\beta_1\left(\beta_2 - \frac{\beta_1^2}{2}\right) + \frac{\varphi_4(x)}{8}\beta_1^3. \quad (3.19)$$

From (3.14) and (3.17)

$$\alpha_1 = -\beta_1, \alpha_1^2 = \beta_1^2 \text{ and } \alpha_1^3 = -\beta_1^3. \quad (3.20)$$

Squaring and adding (3.14) and (3.17)

$$2(\rho + 1)^2 (\Gamma_2[\mu_1; \nu_1])^2 b_2^2 = \frac{\varphi_2^2(x)}{4} (\alpha_1^2 + \beta_1^2). \quad (3.21)$$

Implies

$$b_2^2 = \frac{\varphi_2^2(x)(\alpha_1^2 + \beta_1^2)}{8(\rho+1)^2(\Gamma_2[\mu_1; \nu_1])^2}, \quad (3.22)$$

Adding (3.15) and (3.18)

$$4(2\rho + 1)(\Gamma_2[\mu_1; \nu_1])^2 b_2^2 = \varphi_2(x)(\alpha_2 + \beta_2) + (\varphi_3(x) - \varphi_2(x))\alpha_1^2. \quad (3.23)$$

Applying (3.20) in (3.22)

$$\alpha_1^2 = \frac{4(\rho+1)^2(\Gamma_2[\mu_1; \nu_1])^2 b_2^2}{\varphi_2^2(x)} \quad (3.24)$$

In (3.23), replacing  $\alpha_1^2$

$$4(2\rho + 1)(\Gamma_2[\mu_1; \nu_1])^2 b_2^2 = \varphi_2(x)(\alpha_2 + \beta_2) + \frac{4(\varphi_3(x) - \varphi_2(x))(\rho+1)^2(\Gamma_2[\mu_1; \nu_1])^2 b_2^2}{\varphi_2^2(x)}.$$

Thus

$$b_2^2 = \frac{\varphi_2^3(x)(\alpha_2 + \beta_2)}{4(\Gamma_2[\mu_1; \nu_1])^2[(2\rho+1)\varphi_2^2(x) - (\rho+1)^2(\varphi_3(x) - \varphi_2(x))]} \quad (3.25)$$

Applying Lemma 2.2 and using (1.10)

$$|b_2| \leq \frac{|\delta x| \sqrt{|\delta x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{[(2\rho+1)(\delta x+1) + \rho^2] \delta x - [l\delta x^2 + m\nu] (\rho+1)^2}}.$$

Subtracting (3.18) from (3.15)

$$b_3 = \frac{\varphi_2(x)(\alpha_2 - \beta_2)}{4(2\rho+1)\Gamma_3[\mu_1; \nu_1]} + \frac{(\Gamma_2[\mu_1; \nu_1])^2 b_2^2}{\Gamma_3[\mu_1; \nu_1]}. \quad (3.26)$$

$$b_3 = \frac{\varphi_2(x)(\alpha_2 - \beta_2)}{4(2\rho+1)\Gamma_3[\mu_1; \nu_1]} + \frac{\varphi_2^2(x)\alpha_1^2}{4(\rho+1)^2\Gamma_3[\mu_1; \nu_1]}. \quad (3.27)$$

Applying Lemma 2.2 and using (1.10)

$$|b_3| \leq \frac{|\delta x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{\delta x}{(\rho+1)^2} \right|.$$

By removing (3.19) from (3.16)

$$\begin{aligned} b_4 &= \frac{\varphi_2(x)(\alpha_3 - \beta_3)}{4(3\rho+1)\Gamma_4[\mu_1; \nu_1]} + \frac{[\varphi_3(x) - \varphi_2(x)]\alpha_1(\alpha_2 + \beta_2)}{4(3\rho+1)\Gamma_4[\mu_1; \nu_1]} \\ &+ \frac{5\varphi_2^2(x)\alpha_1(\alpha_2 - \beta_2)}{16(\rho+1)(2\rho+1)\Gamma_4[\mu_1; \nu_1]} + \frac{[\varphi_2(x) - 2\varphi_3(x) + \varphi_4(x)]\alpha_1^3}{8(3\rho+1)\Gamma_4[\mu_1; \nu_1]}. \end{aligned} \quad (3.28)$$

Applying Lemma 2.2 and using (1.10)

$$|b_4| \leq \frac{|l^2 \delta x^3 + l m x \nu + m \delta x|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{5\delta^2 x^2}{2|(\rho+1)(2\rho+1)|\Gamma_4[\mu_1; \nu_1]}$$

The proof of theorem 3.1 is completed.

The subsequent statements are merely corollaries associated with the specific instances of Horadam polynomials.

**Corollary 3.2** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, \mathbb{F}_n(x)]$ . Then

$$|b_2| \leq \frac{|x|\sqrt{|x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(2-\rho)x\rho x + (1+\rho^2)x - (\rho+1)^2|}}$$

$$|b_3| \leq \frac{|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{|x^3+2x|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{5x^2}{2|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

**Corollary 3.3** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, L_n(x)]$ . Then

$$|b_2| \leq \frac{|x|\sqrt{|x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(2-\rho)x\rho x + (1+\rho^2)x - 2(\rho+1)^2|}}$$

$$|b_3| \leq \frac{|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{|x^3+3x|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{5x^2}{2|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

**Corollary 3.4** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, P_n(x)]$ . Then

$$|b_2| \leq \frac{2|x|\sqrt{|2x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(1-\rho)x4\rho x + (1+\rho^2)2x - (\rho+1)^2|}}$$

$$|b_3| \leq \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{2x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{4x|2x^2+1|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{10x^2}{|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

**Corollary 3.5** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, Q_n(x)]$ . Then

$$|b_2| \leq \frac{2|x|\sqrt{|2x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(1-\rho)x4\rho x + (1+\rho^2)2x - 2(\rho+1)^2|}}$$

$$|b_3| \leq \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{2x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{2x|4x^2+3|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{10x^2}{|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

**Corollary 3.6** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, T_n(x)]$ . Then

$$|b_2| \leq \frac{|x|\sqrt{|x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(1-\rho)x2\rho x + (1+\rho^2)x - (1+2\rho)x^2 + (\rho+1)^2|}}$$

$$|b_3| \leq \frac{|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{x|4x^2-3|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{5x^2}{2|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

**Corollary 3.7** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, U_n(x)]$ . Then

$$|b_2| \leq \frac{2|x|\sqrt{|2x|}}{\Gamma_2[\mu_1; \nu_1] \sqrt{|(1-\rho)x4\rho x + (1+\rho^2)2x + (\rho+1)^2|}}$$

$$|b_3| \leq \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]} \left| \frac{1}{(2\rho+1)} + \frac{2x}{(\rho+1)^2} \right|.$$

$$|b_4| \leq \frac{4x|2x^2-1|}{|3\rho+1|\Gamma_4[\mu_1; \nu_1]} + \frac{10x^2}{|\rho+1||2\rho+1|\Gamma_4[\mu_1; \nu_1]}$$

#### 4 Fekete-Szegő Inequality for the class $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, \Phi(x)]$

**Theorem 4.1** If  $\psi$  given by (1.1) is in  $\mathfrak{I}_\sigma^{q,s}[\mu_1; \nu_1, \Phi(x)]$ . Then for some  $\eta \in \mathbb{R}$

$$|b_3 - \eta b_2^2| \leq \begin{cases} 4|\delta x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{|\delta x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

where

$$|\Delta| = \frac{\delta^2 x^2}{4|[(2\rho+1)(\delta x+1)+\rho^2]\delta x-(\rho+1)^2(\delta x^2+m\gamma)|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|. \quad (4.1)$$

**Proof.** From (3.26)

$$b_3 - \eta b_2^2 = \frac{\varphi_2(x)(\alpha_2 - \beta_2)}{4(2\rho+1)\Gamma_3[\mu_1; \nu_1]} + \left[ \frac{(\Gamma_2[\mu_1; \nu_1])^2}{\Gamma_3[\mu_1; \nu_1]} - \eta \right] b_2^2.$$

Using (3.25)

$$b_3 - \eta b_2^2 = \varphi_2(x) \left\{ \left( \Delta(\eta, \rho) + \frac{1}{4(2\rho+1)\Gamma_3[\mu_1; \nu_1]} \right) \alpha_2 + \left( \Delta(\eta, \rho) - \frac{1}{4(2\rho+1)\Gamma_3[\mu_1; \nu_1]} \right) \beta_2 \right\}$$

Where

$$\Delta(\eta, \rho) = \frac{\varphi_2^2(x)}{4[(2\rho+1)\varphi_2^2(x) - (\rho+1)^2(\varphi_3(x) - \varphi_2(x))]} \left[ \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right]$$

Using Lemma 2.1

$$|b_3 - \eta b_2^2| \leq \begin{cases} 4|\delta x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{|\delta x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

**Corollary 4.2** If  $\psi$  given by (1.1) is in  $\mathfrak{I}_\sigma^{q,s}[\mu_1; \nu_1, \mathbb{F}_n(x)]$ . Then

$$|b_3 - \eta b_2^2| \leq \begin{cases} 4|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{4|[\rho x(2-\rho x)+x(1+\rho^2)-(\rho+1)^2]|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

**Corollary 4.3** If  $\psi$  given by (1.1) is in  $\mathfrak{I}_\sigma^{q,s}[\mu_1; \nu_1, L_n(x)]$ . Then

$$|b_3 - \eta b_2^2| \leq \begin{cases} 4|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{4|[\rho x(2-\rho x)+x(1+\rho^2)-2(\rho+1)^2]|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

**Corollary 4.4** If  $\psi$  given by (1.1) is in  $\mathfrak{I}_\sigma^{q,s}[\mu_1; \nu_1, P_n(x)]$ . Then



$$|b_3 - \eta b_2^2| \leq \begin{cases} 8|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{|[4\rho x(1-\rho x)+2x(1+\rho^2)-(\rho+1)^2]|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

**Corollary 4.5** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, Q_n(x)]$ . Then

$$|b_3 - \eta b_2^2| \leq \begin{cases} 8|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{|[4\rho x(1-\rho x)+2x(1+\rho^2)-2(\rho+1)^2]|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

**Corollary 4.6** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, T_n(x)]$ . Then

$$|b_3 - \eta b_2^2| \leq \begin{cases} 4|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{4|[(2\rho+1)(1-x)+\rho^2(1-2x)]x+(\rho+1)^2|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

**Corollary 4.7** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, U_n(x)]$ . Then

$$|b_3 - \eta b_2^2| \leq \begin{cases} 8|x||\Delta| & \text{if } |\Delta| \geq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \\ \frac{2|x|}{\Gamma_3[\mu_1; \nu_1]|2\rho+1|} & \text{if } |\Delta| \leq \frac{1}{4\Gamma_3[\mu_1; \nu_1]|2\rho+1|} \end{cases}$$

Where

$$|\Delta| = \frac{x^2}{|[4\rho x(1-\rho x)+2x(1+\rho^2)+(\rho+1)^2]|} \left| \frac{1}{\Gamma_3[\mu_1; \nu_1]} - \frac{\eta}{(\Gamma_2[\mu_1; \nu_1])^2} \right|.$$

### 5 Second Hankel Determinant for $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, \Phi(x)]$

**Theorem 5.1** If  $\psi$  given by (1.1) is in  $\mathfrak{F}_\sigma^{q,s}[\mu_1; \nu_1, \Phi(x)]$ . Then

$$|b_2 b_4 - b_3^2| \leq \begin{cases} \mathcal{J}(x, 2); & B_1 \geq 0 \text{ and } B_2 \geq 0 \\ \max \left\{ \frac{\delta^2 x^2}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}, \mathcal{J}(x, 2) \right\}; & B_1 > 0 \text{ and } B_2 < 0 \\ \frac{\delta^2 x^2}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}; & B_1 \leq 0 \text{ and } B_2 \leq 0 \\ \max \{ \mathcal{J}(x, \alpha_0), \mathcal{J}(x, 2) \}; & B_1 < 0 \text{ and } B_2 > 0 \end{cases}$$

Where

$$\mathcal{J}(x, 2) = \frac{\delta^2 x^2}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} + \frac{B_1 + B_2}{2(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

$$\mathcal{J}(x, \alpha_0) = \frac{\delta^2 x^2}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} - \frac{B_2^2}{8B_1(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

$$\begin{aligned}
B_1 &= 2(\rho + 1)^3(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2 \delta x[(l\delta x^2 + m\gamma)(lx - 2) + \delta x(m - 1)] \\
&+ 2(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]\delta^2 x^2[(\rho + 1)^4 - (2\rho + 1)^2\delta^2 x^2] \\
&+ (5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1])(\rho + 1)^2(2\rho + 1)(3\rho + 1)\delta^3 x^3,
\end{aligned}$$

and

$$\begin{aligned}
B_2 &= 2(\rho + 1)^3(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2 \delta x[\delta x(1 + 2lx) + 2m\gamma] \\
&- 4(\rho + 1)^4(3\rho + 1)\delta^2 x^2\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1] \\
&+ (5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1])(\rho + 1)^2(2\rho + 1)(3\rho + 1)\delta^3 x^3.
\end{aligned}$$

**Proof.** From (3.14), (3.27) and (3.28)

$$\begin{aligned}
b_2 b_4 - b_3^2 &= \frac{\varphi_2(x)[\varphi_2(x) - 2\varphi_3(x) + \varphi_4(x)](\rho + 1)^3(\Gamma_3[\mu_1; \nu_1])^2 - \varphi_2^4(x)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}{16(\rho + 1)^4(3\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \alpha_1^4 \\
&+ \frac{\varphi_2(x)[\varphi_3(x) - \varphi_2(x)]\alpha_1^2(\alpha_2 + \beta_2)}{8(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} + \frac{\varphi_2^2(x)\alpha_1(\alpha_3 - \beta_3)}{8(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} \\
&+ \frac{[5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]]\varphi_2^3(x)\alpha_1^2(\alpha_2 - \beta_2)}{32(\rho + 1)^2(2\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} - \frac{\varphi_2^2(x)(\alpha_2 - \beta_2)^2}{16(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2}.
\end{aligned}$$

using Lemma[2.3]

$$\alpha_2 - \beta_2 = \frac{(4 - \alpha_1^2)(x - y)}{2}, \quad (5.1)$$

$$\alpha_2 + \beta_2 = \frac{(4 - \alpha_1^2)(x + y)}{2} + \alpha_1^2, \quad (5.2)$$

and

$$\alpha_3 - \beta_3 = \frac{\alpha_1^3}{2} + \frac{4 - \alpha_1^2}{2} \alpha_1(x + y) - \frac{4 - \alpha_1^2}{4} \alpha_1(x^2 + y^2) + \frac{4 - \alpha_1^2}{2} [(1 - |x|^2)z - (1 - |y|^2)w]. \quad (5.3)$$

for some  $x, y, z, w$  with  $|x| \leq 1$ ,  $|y| \leq 1$ ,  $|z| \leq 1$ ,  $|w| \leq 1$ ,  $|\alpha_1| \in (0, 2)$  and substituting  $\alpha_2 - \beta_2$ ,  $\alpha_2 + \beta_2$  and  $\alpha_3 - \beta_3$  and after some simplifications

$$\begin{aligned}
b_2 b_4 - b_3^2 &= \frac{\varphi_2(x)\varphi_4(x)(\rho + 1)^3(\Gamma_3[\mu_1; \nu_1])^2 - \varphi_2^4(x)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}{16(\rho + 1)^4(3\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \alpha_1^4 \\
&+ \frac{\varphi_2(x)\varphi_3(x)\alpha_1^2(4 - \alpha_1^2)(x + y)}{16(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} - \frac{\varphi_2^2(x)(4 - \alpha_1^2)\alpha_1^2(x^2 + y^2)}{32(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} \\
&+ \frac{[5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]]\varphi_2^3(x)\alpha_1^2(4 - \alpha_1^2)(x - y)}{64(\rho + 1)^2(2\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} - \frac{\varphi_2^2(x)(4 - \alpha_1^2)^2(x - y)^2}{64(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2} \\
&+ \frac{\varphi_2^2(x)(4 - \alpha_1^2)\alpha_1[(1 - |x|^2)z - (1 - |y|^2)w]}{16(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}.
\end{aligned}$$

Let  $\alpha = \alpha_1$ , without any restriction it can be assumed that  $\alpha \in [0, 2]$ ,  $\xi_1 = |x| \leq 1$ ,  $\xi_2 = |y| \leq 1$  and applying triangular inequality,

$$\begin{aligned}
|b_2 b_4 - b_3^2| &\leq \frac{\varphi_2(x)\varphi_4(x)(\rho + 1)^3(\Gamma_3[\mu_1; \nu_1])^2 - \varphi_2^4(x)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}{16(\rho + 1)^4(3\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \alpha^4 \\
&+ \frac{\varphi_2^2(x)\alpha(4 - \alpha^2)}{8(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} + \frac{\varphi_2^2(x)(4 - \alpha^2)\alpha(\alpha - 2)(\xi_1^2 + \xi_2^2)}{32(\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} \\
&+ \frac{4(\rho + 1)(2\rho + 1)(\Gamma_3[\mu_1; \nu_1])^2\varphi_2(x)\varphi_3(x)\alpha^2(4 - \alpha^2)(\xi_1 + \xi_2)}{64(\rho + 1)^2(2\rho + 1)(3\rho + 1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}
\end{aligned}$$

$$+ \frac{(5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1])\varphi_2^3(x)(3\rho+1)\alpha^2(4-\alpha^2)(\xi_1+\xi_2)}{64(\rho+1)^2(2\rho+1)(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

$$+ \frac{\varphi_2^2(x)(4-\alpha^2)^2(\xi_1+\xi_2)^2}{64(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}.$$

That is

$$|b_2b_4 - b_3^2| \leq T_1(x, \alpha) + T_2(x, \alpha)(\xi_1 + \xi_2) + T_3(x, \alpha)(\xi_1^2 + \xi_2^2) + T_4(x, \alpha)(\xi_1 + \xi_2)^2$$

$$= F(\xi_1, \xi_2). \quad (5.4)$$

Where,

$$T_1(x, \alpha) = \frac{\varphi_2(x)\varphi_4(x)(\rho+1)^3(\Gamma_3[\mu_1; \nu_1])^2 - \varphi_2^4(x)(3\rho+1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}{16(\rho+1)^4(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \alpha^4$$

$$+ \frac{\varphi_2^2(x)\alpha(4-\alpha^2)}{8(\rho+1)(3\rho+1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} \geq 0,$$

$$T_2(x, \alpha) = \frac{4(\rho+1)(2\rho+1)(\Gamma_3[\mu_1; \nu_1])^2\varphi_2(x)\varphi_3(x)\alpha^2(4-\alpha^2)}{64(\rho+1)^2(2\rho+1)(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

$$+ \frac{(5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1])\varphi_2^3(x)(3\rho+1)\alpha^2(4-\alpha^2)}{64(\rho+1)^2(2\rho+1)(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \geq 0,$$

$$T_3(x, \alpha) = \frac{\varphi_2^2(x)(4-\alpha^2)\alpha(\alpha-2)}{32(\rho+1)(3\rho+1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]} \leq 0,$$

$$T_4(x, \alpha) = \frac{\varphi_2^2(x)(4-\alpha^2)^2}{64(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} \geq 0, \quad \text{and} \quad 0 \leq \alpha \leq 2.$$

As we have to maximize  $F(\xi_1, \xi_2)$  in the closed square

$$C_s = \{(\xi_1, \xi_2), 0 \leq \xi_1 \leq 1, 0 \leq \xi_2 \leq 1\}.$$

Let  $\alpha = 0$ ,  $\alpha = 2$  and  $\alpha \in (0, 2)$

When  $\alpha = 0$ ,

$$F(\xi_1, \xi_2) = T_4(x, 0) = \frac{\varphi_2^2(x)(\xi_1+\xi_2)^2}{4(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}$$

and in this case  $F(\xi_1, \xi_2)$  reaches it maximum at  $(\xi_1, \xi_2)$  and

$$\max\{F(\xi_1, \xi_2): \xi_1, \xi_2 \in [0, 1]\} = F(1, 1) = \frac{\varphi_2^2(x)}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}.$$

When  $\alpha = 2$ ,  $F(\xi_1, \xi_2)$  is a constant function.

$$F(\xi_1, \xi_2) = T_1(x, 2) = \frac{\varphi_2(x)\varphi_4(x)(\rho+1)^3(\Gamma_3[\mu_1; \nu_1])^2 - \varphi_2^4(x)(3\rho+1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]}{(\rho+1)^4(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}.$$

When  $\alpha \in (0, 2)$ , let  $\xi_1 + \xi_2 = e$  and  $\xi_1 \cdot \xi_2 = g$ , then

$$F(\xi_1, \xi_2) = T_1(x, \alpha) + T_2(x, \alpha)e + (T_3(x, \alpha) + T_4(x, \alpha))e^2 - 2T_3(x, \alpha)g = H(e, g).$$

where  $e \in [0, 2]$  and  $g \in [0, 1]$ . Now we have to investigate the maximum of

$$H(e, g) \in \Theta = \{(e, g): e \in [0, 2], g \in [0, 1]\} \quad (5.5)$$

Differentiating 5.5 partially,

$$\frac{\partial H}{\partial e} = T_2(x, \alpha) + 2(T_3(x, \alpha) + T_4(x, \alpha))e = 0 \quad \text{and} \quad \frac{\partial H}{\partial g} = -2T_3(x, \alpha) = 0$$

reveals  $H(e, g)$  does not have a critical point in closed square  $C_s$  and so  $F(\xi_1, \xi_2)$  does not have a critical point in the square  $C_s$ .

So  $F(\xi_1, \xi_2)$  cannot have maximum value in the interior of closed square. The maximum of  $F(\xi_1, \xi_2)$  on the

boundary of  $C_s$  has to be investigated.

For  $\xi_1 = 0, \xi_2 \in [0,1]$  (similarly for  $\xi_2 = 0, \xi_1 \in [0,1]$ ) and

$$F(0, \xi_2) = T_1(x, \alpha) + T_2(x, \alpha)\xi_2 + [T_3(x, \alpha) + T_4(x, \alpha)]\xi_2^2 = M(\xi_2).$$

Since  $T_3(x, \alpha) + T_4(x, \alpha) \geq 0$ , we have

$$M'(\xi_2) = T_2(x, \alpha) + 2[T_3(x, \alpha) + T_4(x, \alpha)]\xi_2 > 0,$$

which implies that  $M(\xi_2)$  is an increasing function. Therefore for a fixed  $\alpha \in [0,2)$  the maximum occurs at  $\xi_2 = 1$  and

$$\text{Max } M(\xi_2) = M(1) = T_1(x, \alpha) + T_2(x, \alpha) + T_3(x, \alpha) + T_4(x, \alpha) = F(0,1).$$

For  $\xi_1 = 1, \xi_2 \in [0,1]$  (similarly for  $\xi_2 = 1, \xi_1 \in [0,1]$ ) and

$$F(1, \xi_2) = T_1(x, \alpha) + T_2(x, \alpha) + T_3(x, \alpha) + T_4(x, \alpha) + [T_2(x, \alpha) + 2T_4(x, \alpha)]\xi_2 + [T_3(x, \alpha) + T_4(x, \alpha)]\xi_2^2 = G(\xi_2).$$

As  $T_3(x, \alpha) + T_4(x, \alpha) \geq 0$  then

$$G'(\xi_2) = T_2(x, \alpha) + 2T_4(x, \alpha) + 2[T_3(x, \alpha) + T_4(x, \alpha)]\xi_2 > 0.$$

Therefore  $G(\xi_2)$  is an increasing function and the maximum is attained at  $\xi_2 = 1$ .

$$\text{Max}\{F(1, \xi_2); \xi_2 \in [0,1]\} = F(1,1) = T_1(x, \alpha) + 2[T_2(x, \alpha) + T_3(x, \alpha)] + 4T_4(x, \alpha). \quad (5.6)$$

Now for each  $\alpha \in (0,2)$ , we have

$$T_1(x, \alpha) + 2[T_2(x, \alpha) + T_3(x, \alpha)] + 4T_4(x, \alpha) > T_1(x, \alpha) + T_2(x, \alpha) + T_3(x, \alpha) + T_4(x, \alpha).$$

Thus

$\max\{F(\xi_1, \xi_2); \xi_1 \in [0,1], \xi_2 \in [0,1]\} = T_1(x, \alpha) + 2[T_2(x, \alpha) + T_3(x, \alpha)] + 4T_4(x, \alpha)$ . Since  $M(1) \leq G(1)$  for  $\alpha \in [0,2]$ . we have,  $\max\{F(\xi_1, \xi_2)\} = F(1,1)$  occurs on the boundary of closed square.

Let  $J: (0,2) \rightarrow \mathbb{R}$  be the function defined by

$$J(x, \alpha) = \max\{F(\xi_1, \xi_2)\} = F(1,1) = T_1(x, \alpha) + 2[T_2(x, \alpha) + T_3(x, \alpha)] + 4T_4(x, \alpha). \quad (5.7)$$

By putting the values of  $T_1(x, \alpha), T_2(x, \alpha), T_3(x, \alpha), T_4(x, \alpha)$  in (5.7)

$$J(x, \alpha) = \frac{\varphi_2^2(x)}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} + \frac{B_1\alpha^4 + 4B_2\alpha^2}{32(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} \quad (5.8)$$

where

$$B_1 = 2(\rho + 1)^3(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2\varphi_2(x)[\varphi_4(x) - 2\varphi_3(x) - \varphi_2(x)] + 2(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]\varphi_2^2(x)[(\rho + 1)^4 - (2\rho + 1)^2\varphi_2^2(x)] + [5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]](\rho + 1)^2(2\rho + 1)(3\rho + 1)\varphi_2^3(x).$$

and

$$B_2 = 2(\rho + 1)^3(2\rho + 1)^2(\Gamma_3[\mu_1; \nu_1])^2\varphi_2(x)[2\varphi_3(x) + \varphi_2(x)] - 4(\rho + 1)^4(3\rho + 1)\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]\varphi_2^2(x) + [5(\Gamma_3[\mu_1; \nu_1])^2 - 4\Gamma_2[\mu_1; \nu_1]\Gamma_4[\mu_1; \nu_1]](\rho + 1)^2(2\rho + 1)(3\rho + 1)\varphi_2^3(x).$$

If  $J(x, \alpha)$  has a maximum value in the interior of  $\alpha \in [0,2]$  we have

$$J'(x, \alpha) = \frac{B_1\alpha^3 + 2B_2\alpha}{8(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

As the sign of  $J'(x, \alpha)$  depends on the sign of  $B_1$  and  $B_2$ , we have the following cases

**Result 1:** When  $B_1 \geq 0$  and  $B_2 \geq 0$  we have  $J'(x, \alpha) \geq 0$ . So  $J(x, \alpha)$  is an increasing function. Therefore

$$\max\{J(x, \alpha); \alpha \in (0, 2)\} = J(x, 2)$$

$$= \frac{\varphi_2^2(x)}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} + \frac{B_1+B_2}{2(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

Thus  $\text{Max}\{\text{Max}\{F(\xi_1, \xi_2); 0 \leq \xi_1, \xi_2 \leq 1\}; 0 < \alpha < 2\} = J(x, 2)$

**Result 2:** When  $B_1 > 0$  and  $B_2 < 0$ . we have,

$$J'(x, \alpha) = \frac{(B_1\alpha^2+2B_2)\alpha}{8(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

So  $J'(x, \alpha) = 0$  at point  $\alpha_0 = \sqrt{\frac{-2B_2}{B_1}}$  which is the critical point of  $J(x, \alpha)$ .

Now

$$J''(x, \alpha) = \frac{-4B_2}{8(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]} > 0 \text{ at } \alpha_0.$$

Therefore  $\alpha_0$  is the minimum point of the function  $J(x, \alpha)$ . Hence  $J(x, \alpha)$  cannot have a maximum.

**Result 3:** If  $B_1 \leq 0$  and  $B_2 \leq 0$  then  $J'(x, \alpha) \leq 0$ , therefore  $J(x, \alpha)$  is a decreasing function on the interval  $(0, 2)$ . Hence

$$\max\{J(x, \alpha); \alpha \in (0, 2)\} = J(x, 0) = \frac{\varphi_2^2(x)}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2}. \quad (5.9)$$

**Result 4:** If  $B_1 < 0$  and  $B_2 > 0$  then  $J''(x, \alpha) < 0$  at  $\alpha_0$ . Hence  $\alpha_0$  is the maximum point of  $J(x, \alpha)$  and the maximum value occurs at  $\alpha = \alpha_0$ . Thus

$$\max\{J(x, \alpha); \alpha \in (0, 2)\} = J(x, \alpha_0)$$

$$J(x, \alpha_0) = \frac{\varphi_2^2(x)}{(2\rho+1)^2(\Gamma_3[\mu_1; \nu_1])^2} - \frac{B_2^2}{8B_1(\rho+1)^4(2\rho+1)^2(3\rho+1)\Gamma_2[\mu_1; \nu_1](\Gamma_3[\mu_1; \nu_1])^2\Gamma_4[\mu_1; \nu_1]}$$

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